

## Brief Summary of Nows Theoretical Backgrounds

### ☉ Power spectral density matrix of longitudinal wind velocity fluctuation

- ▣ Power spectral density function of longitudinal wind velocity fluctuations – two-sided form
- Kaimal et al. (1972); Simiu (1974); Simiu and Scanlan (1996),

$$S_{rr}(z, \omega) = \frac{1}{2} \cdot \frac{200}{2\pi} \cdot u_*^2 \cdot \frac{z}{U(z)} \cdot \frac{1}{\left[1 + 50 \cdot \frac{\omega z}{2\pi U(z)}\right]^{5/3}}$$

where,  $z$  = height,  $\omega$  = circular frequency (rad/s);  $u_*$  = friction velocity;  $U(z)$  = mean wind speed at height  $z$

- ▣ Coherence function (two-dimensional)
- Davenport (1967); Simiu and Scanlan (1996)

$$f_{rs}(\omega) = \exp\left[-\frac{\omega}{2\pi} \frac{C_z \Delta z}{\frac{1}{2}[U(z_r) + U(z_s)]}\right] \cdot \exp\left[-\frac{\omega}{2\pi} \frac{C_x \Delta x}{\frac{1}{2}[U(x_r) + U(x_s)]}\right]$$

where,  $x, z$  = horizontal and vertical directions, respectively;  $\Delta z = |z_r - z_s|$ ;  $\Delta x = |x_r - x_s|$ ;  $C_z,$

$C_x$  = a constant, generally taken 10 and 16 for structural design viewpoint, respectively.

- ▣ Cross-spectral density function
- Co-spectrum (quadratic term of wind is ignored)

$$S_{rs}(\omega) = \sqrt{S_{rr}(\omega) \cdot S_{ss}(\omega)} \exp(-f_{rs}(\omega))$$

- ▣ Power spectral density matrix  $S(\omega)$  : two-dimensional,  $n$ -variate

$$S(\omega) = \begin{bmatrix} S_{11}(\omega) & S_{12}(\omega) & \cdots & S_{1n}(\omega) \\ S_{21}(\omega) & S_{22}(\omega) & \cdots & S_{2n}(\omega) \\ \vdots & \vdots & \cdots & \vdots \\ S_{n1}(\omega) & S_{n2}(\omega) & \cdots & S_{nn}(\omega) \end{bmatrix}$$

## © Simulation schemes of wind velocity fluctuations

### 1. Discrete frequency function with FFT

– Wittig and Sinha (1975)

Discrete time series can be simulated using the following model:

$$y_p(n\Delta t) = \frac{1}{N} \sum_{k=0}^N Y_p(k\Delta f) \exp\left(j \frac{2\pi kn}{N}\right)$$

$$Y_p(k\Delta f) = \sum_{i=0}^p H_{pi}(k\Delta f) \varepsilon_{ik} \sqrt{2f_c N}$$

where,

$H_{pi}(k\Delta f)$  : a lower triangular matrix by Cholesky decomposition of one-sided power spectral density function  $S(f)$ ;  $\varepsilon_{ik} = \xi_{ik} + j\eta_{ik}$  = complex Gaussian random number with zero mean and

0.5 variance;  $\Delta t = \frac{1}{2f_c}$ ;  $f_c$  = Nyquist frequency

## 2. Schur decomposition approach with AR (autoregressive)

– Di Paola (1998); Di Paola and Gullo (2001)

The  $n$ -variate stochastic vector process  $V(t)$  can be decomposed into a summation of  $n$ -variate fully coherent normal vectors  $Y_j(t)$  independent of each other:

$$V(t) = \sum_{j=1}^n Y_j(t)$$

Let  $\psi(\omega)$  be the eigenmatrix of  $S(\omega)$  whose columns are the eigenvectors (real and orthogonal), then, following relationship holds:

$$\Psi^T(\omega)S(\omega)\Psi(\omega) = \Lambda(\omega)$$

$$\Psi^T(\omega)\Psi(\omega) = I$$

Vectors  $Y_j(t)$  can be described as:

$$Y_j(t) = \int_{-\infty}^{\infty} S(\omega)e^{i\omega t} dB_j(\omega) = \int_{-\infty}^{\infty} \psi_j(\omega)\sqrt{\Lambda(\omega)}e^{i\omega t} dB_j(\omega)$$

Let define the frequency domain  $[\omega_0, \omega_c]$ , where  $\omega_0$  and  $\omega_c$  are lower and upper cut-off frequencies, and subdivided the domain into  $M$  parts  $\omega_0 \equiv \Omega_0, \Omega_1, \dots, \Omega_m \equiv \omega_c$

With third-order polynomial approximation of eigenvectors  $\psi_j^{(s)}(\omega)$ ,

$$\psi_j^{(s)}(\omega) = N_j^{(s)}l(\omega), \quad \Omega_{s-1} \leq \omega \leq \Omega_s$$

where,  $l(\omega)=[1 \ \omega \ \omega^2 \ \omega^3]$ .

Accordingly, vectors  $Y_j(t)$  can be expressed as:

$$Y_j(t) = \sum_{s=1}^M N_j^{(s)} \int_{\Omega_{s-1}}^{\Omega_s} l(\omega)\sqrt{\Lambda(\omega)}e^{i\omega t} dB_j^{(s)}(\omega) = \sum_{s=1}^M N_j^{(s)}U_j^{(s)}(t)$$

$$\text{where, } U_j^{(s)}(t) = \sum_{s=1}^M \int_{\Omega_{s-1}}^{\Omega_s} l(\omega)\sqrt{\Lambda(\omega)}e^{i\omega t} dB_j^{(s)}(\omega)$$

Using the standard generation via AR(autoregressive) model:

$$U_{j,r}^{(s)}(t_k) = \sum_{u=1}^p a_{j,u}^{(s)}U_{j,r}^{(s)}(t_{k-u}) + \sigma_{j,1}^{(s)}W_j^{(s)}(t_k)$$

where  $r = 1, \dots, 4$ ;  $a_{j,u}^{(s)}$  = parameters of the AR model;  $\sigma_{j,r}^{(s)}$  = variances of the input;  $W_j^{(s)}$  = normal random variables with zero mean and unit variance;  $p$  = AR model order

AR parameters can be evaluated by using Yule-walker scheme, that is, autocorrelation method.

It is worth noting that in this study,  $M = 1$  and  $p = 4$  are used to generate wind velocity fluctuations.

### 3. Ergodic spectral representation method

– Deodatis (1996); Ding et al. (2006)

Power spectral density matrix  $S(\omega)$  can be decomposed into the following product:

$$S(\omega) = H(\omega)H^{T*}(\omega)$$

where,  $H(\omega)$  is a lower triangular matrix by Cholesky decomposition of  $S(\omega)$

Stochastic process  $V_j(t)$  can be described by using following trigonometric series:

$$V_j(t) = 2 \sum_{m=1}^j \sum_{l=1}^N |H_{jm}(\omega)| \sqrt{\Delta\omega} \cos[\omega_{ml}t - \theta_{jm}(\omega_{ml}) + \Phi_{ml}]$$

To take advantage of the FFT technique, above equation can be rewritten as follows:

$$V_j(p\Delta t) = \text{Re} \left\{ \sum_{m=1}^j h_{jm}(p\Delta t) \exp \left[ i \left( \frac{m\Delta\omega}{n} \right) (p\Delta t) \right] \right\}$$

where,  $N$  = number of n-variate simulation;  $j = 1, 2, \dots, N$ ;  $p = 0, 1, \dots, n \times (M - 1)$ ;  $M = 2N$ ;

$$h_{jm}(p\Delta t) = \begin{cases} g_{jm}(p\Delta t) & \text{for } p = M, M+1, \dots, 2M-1 \\ g_{jm}[(p-M)\Delta t] & \text{for } p = M, M+1, \dots, 2M-1 \\ \vdots & \vdots \\ g_{jm}[(p-nM)\Delta t] & \text{for } p = nM, n(M+1), \dots, nM-1 \end{cases}$$

where,

$$g_{jm}(p\Delta t) = \sum_{l=0}^{M-1} B_{jml} \exp \left[ ilp \frac{2\pi}{M} \right] \quad p = 0, 1, \dots, M-1$$

$$B_{jml} = 2 |H_{jm}(\omega_{ml})| \sqrt{\Delta\omega} \exp[-i\theta_{jm}(\omega_{ml})] \exp(i\Phi_{ml})$$

$$\omega_{ml} = (l-1)\Delta\omega + \frac{m}{n}\Delta\omega \quad m = 1, 2, \dots, n; \quad l = 1, 2, \dots, N$$

$$\theta_{jm}(\omega_{ml}) = \tan^{-1} \left\{ \frac{\text{Im}(H(\omega_{ml}))}{\text{Re}(H(\omega_{ml}))} \right\}$$

$\Phi_{ml}$  = independent random phase angles distributed uniformly over the interval  $[0, 2\pi]$

Note :  $g_{jm}(p\Delta t)$  can be obtained from inverse FFT of  $B_{jml}$

#### 4. Conventional spectral representation method

– Shinozuka and Deodatis (1991)

Power spectral density matrix  $S(\omega)$  can be decomposed into the following product:

$$S(\omega) = H(\omega)H^{T*}(\omega)$$

where,  $H(\omega)$  is a lower triangular matrix by Cholesky decomposition of  $S(\omega)$ .

Stochastic process  $V(t)$  can be described by using following trigonometric series:

$$V(t) = \sqrt{2} \sum_{n=0}^{N-1} A_n \cos(\omega_n t + \phi_n)$$

To take advantage of the FFT technique, above equation can be rewritten as follows:

$$V(p\Delta t) = \text{Re} \left\{ \sum_{n=0}^{M-1} B_n \exp[i(n\Delta\omega)(p\Delta t)] \right\}$$

where,  $A_n = \sqrt{(2 \cdot \Delta\omega)} |H(\omega)|$ ,  $N =$  number of  $n$ -variate simulation;  $p = 0, 1, \dots, M - 1$ ;

$M = 2N$ ;  $B_n = 2 |H(\omega)| \sqrt{\Delta\omega} \cdot \exp(i\phi_n)$ ;  $\Delta\omega = \omega_u/N$ ;  $\omega_u =$  upper cut-off frequency [rad/sec];

$\phi_n =$  independent random phase angles distributed uniformly over the interval  $[0, 2\pi]$

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